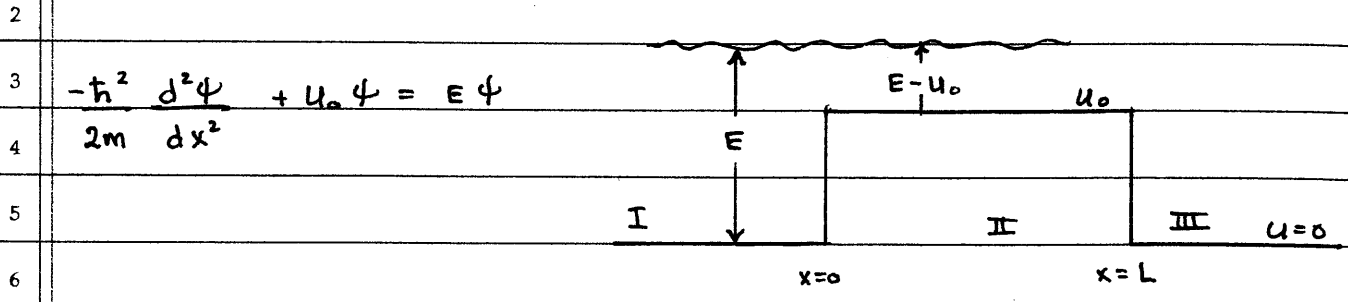


The Schrödinger Equation

1 Constant Potential Energy (a free particle)



3
$$-\hbar^2 \frac{d^2\psi}{2m dx^2} + U_0 \psi = E \psi$$

7
$$\frac{d^2\psi}{dx^2} + (E - U_0) \frac{2m}{\hbar^2} \psi = 0$$

9 Looks like $\frac{d^2\psi}{dx^2} + k^2 \psi = 0$ whose solution $\Rightarrow \psi(x) = A \sin kx + B \cos kx$
 10 and $k = \sqrt{\frac{2m(E - U_0)}{\hbar^2}}$ in region II

12 In regions I and III $k = \sqrt{\frac{2mE}{\hbar^2}}$ because $U_0 = 0$

14 What happens when $E < U_0$?

15 Then, the Schrödinger Equation becomes:

16 Regions I & III $-\hbar^2 \frac{d^2\psi}{2m dx^2} = E \psi \Rightarrow \frac{d^2\psi}{dx^2} + \underbrace{\frac{2mE}{\hbar^2}}_k^2 \psi = 0$

18 $k^2 = \frac{2mE}{\hbar^2}$ $\psi_{I,III} \sim A \cos kx + B \sin kx$
 19 $E \cos kx + F \sin kx$

20 Region II $-\hbar^2 \frac{d^2\psi}{2m dx^2} + U_0 \psi = E \psi \Rightarrow \frac{d^2\psi}{dx^2} - \frac{(U_0 - E) 2m}{\hbar^2} \psi = 0$

22 $\frac{d^2\psi}{dx^2} - K^2 \psi = 0$ $K^2 = \frac{2m(U_0 - E)}{\hbar^2} > 0$

24 $\psi_{II} \sim c e^{-Kx} + d e^{+Kx}$

26 Free Particle Let $U_0 = 0$ $-\hbar^2 \frac{d^2\psi}{2m dx^2} = E \psi \Rightarrow \frac{d^2\psi}{dx^2} + \frac{2mE}{\hbar^2} \psi = 0$

28 $\frac{d^2\psi}{dx^2} + k^2 \psi = 0$ $\psi = A \sin kx + B \cos kx = \frac{A}{2i} (e^{ikx} - e^{-ikx}) + \frac{B}{2} (e^{ikx} + e^{-ikx})$

The Schrödinger Equation

1 Free Particle cont'd

$$2 \quad \psi(x) = \left(\frac{A}{2i} + \frac{B}{2} \right) e^{ikx} + \left(\frac{B}{2} - \frac{A}{2i} \right) e^{-ikx} = A' e^{ikx} + B' e^{-ikx}$$

$$5 \quad \psi(x,t) = \psi(x) e^{-i\omega t} \Rightarrow \text{time-evolution of the wave function}$$

$$7 \quad \psi(x,t) = A' e^{i(kx-\omega t)} + B' e^{-i(kx+\omega t)}$$

8 $(L \rightarrow R) \qquad (R \rightarrow L)$

9 Eigenfunction

$$10 \quad \text{Equation } \text{Pop} \left(A' e^{i(kx-\omega t)} \right) = \frac{\hbar}{i} \frac{d}{dx} \left(A' e^{i(kx-\omega t)} \right) = \hbar (ik) \left(A' e^{i(kx-\omega t)} \right)$$

$$12 \quad \text{Pop} \left(A' e^{i(kx-\omega t)} \right) = \hbar k \left(A' e^{i(kx-\omega t)} \right)$$

13 momentum eigenvalue

$$15 \quad A' \left(e^{i(kx-\omega t)} \right) \Rightarrow A' e^{\frac{i}{\hbar} (\hbar k x - \hbar \omega t)} = A' e^{\frac{i}{\hbar} (p x - E t)}$$

17 The wave can be seen to have

$$\text{momentum} \rightarrow p = \hbar k$$

18 and

$$\text{energy} \rightarrow E = \hbar \omega$$

$$20 \quad P(x) = \psi^*(x) \psi(x)$$

$$21 \quad P(x,t) = \psi^*(x) \psi(x) e^{i\omega t} e^{-i\omega t} = \psi^*(x) \psi(x) e^0 = \underbrace{\psi^*(x) \psi(x)}_{\text{probability/length}}$$

23 Probability density is a constant. (constant in time)

24 This is consistent with the Heisenberg Uncertainty Principle.

$$25 \quad k \rightarrow \text{is precisely defined} \quad \Delta k = 0 \Rightarrow \Delta p = \hbar \Delta k = 0 \quad \Delta x \rightarrow \infty$$

The Schrödinger Equation

Infinite Potential Energy Well

Regions I & III $-\frac{\hbar^2}{2m} \frac{d^2\psi}{dx^2} + \infty\psi = E\psi$

 $u = \infty$ $u = \infty$

I

II

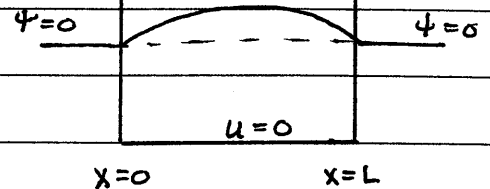
III

$\psi = 0$ is the only solution.

Region II $-\frac{\hbar^2}{2m} \frac{d^2\psi}{dx^2} + 0\psi = E\psi$

$$\psi(x) = A \sin kx + B \cos kx$$

$$\text{where } k^2 = \frac{2mE}{\hbar^2}$$



Apply the Boundary Conditions:

@ $x=0$ $\psi_I(0) = \psi_{II}(0) \Rightarrow 0 = A \sin(0) + B \cos(0) = B$

$$\boxed{\text{So, } B=0}$$

@ $x=L$ $\psi_{II}(L) = \psi_{III}(L) \Rightarrow A \sin(kL) = 0$

This occurs when $k = \frac{n\pi}{L}$ So, $\boxed{\psi_{II}(x) = A \sin k_n x}$

How do you find A? Use the normalization requirement.

$$1 = \int_{-\infty}^{\infty} \psi^*(x) \psi(x) dx = \underbrace{\int_{-\infty}^0 \psi_I^*(x) \psi_I(x) dx}_{=0} + A^2 \underbrace{\int_0^L \sin^2\left(\frac{n\pi x}{L}\right) dx}_{L/2} + \underbrace{\int_L^{\infty} \psi_{III}^*(x) \psi_{III}(x) dx}_{=0}$$

$$A^2 = \frac{2}{L}$$

$$\text{or } A = \sqrt{\frac{2}{L}}$$

$$\boxed{\psi_{II}(x) = \sqrt{\frac{2}{L}} \sin\left(\frac{n\pi x}{L}\right)} = \psi_n(x)$$

$$E_n = \frac{p_n^2}{2m} = \frac{(\hbar k_n)^2}{2m} = \frac{(\hbar \left(\frac{n\pi}{L}\right))^2}{2m}$$

$$\boxed{E_n = \frac{n^2 \pi^2 \hbar^2}{2mL^2}} = n^2 E_1$$

Schrödinger Equation

one dimensional infinitely deep potential well

Example 3.2 Electron trapped in a one-dimensional well of length $L = 0.10 \text{ nm}$.

a.) Find E_1, E_2, E_3

$$E_n = \frac{n^2 \pi^2 \hbar^2}{2m L^2} \Rightarrow E_1 = \frac{\pi^2 (\hbar c)^2}{2m c^2 L^2} = \frac{\pi^2 (197 \text{ eV} \cdot \text{nm})^2}{2(511,000 \text{ eV})(.01 \text{ nm}^2)}$$

$$E_1 = 37.6 \text{ eV}$$

$$E_2 = 150.4 \text{ eV}$$

$$E_3 = 338.4 \text{ eV}$$

b.) How much energy is required to excite the electron from

$$n=1 \rightarrow n=3$$

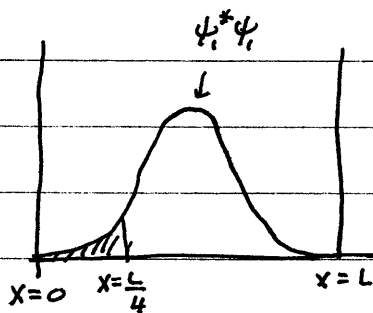
$$\Delta E = E_3 - E_1 = 300.8 \text{ eV}$$

c.) The electron drops down from $n=3 \rightarrow n=2$.

How much energy is released in the process.

~~$$E^* + E_2 = E_3$$~~

$$E^* = E_3 - E_2 = 188.0 \text{ eV}$$



Find the probability of locating a particle between $0 \leq x \leq L/4$ in an infinitely deep potential well. in the $n=1$ state.

$$\frac{2}{L} \int_0^{L/4} \sin^2\left(\frac{\pi x}{L}\right) dx = \frac{2}{L} \int_0^{L/4} \left(1 - \cos \frac{2\pi x}{L}\right) dx$$

$$\text{Prob} = \frac{1}{L} \left[\int_0^{L/4} dx - \int_0^{L/4} \cos \frac{2\pi x}{L} dx \right] = \frac{1}{L} \left[\frac{L}{4} - \frac{L}{2\pi} \sin \frac{2\pi x}{L} \right]_0^{L/4}$$

$$\text{Prob} = \left[\frac{1}{4} - \frac{1}{2\pi} (1-0) \right] = 0.0909$$

$$0.25 - 0.1591 = 0.090845 \text{ (Mathematica)}$$

Schrödinger Equation

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$$\sigma = \sqrt{\overline{x^2} - \bar{x}^2} \quad \bar{x} = \int_0^L x \psi^* \psi dx = \frac{2}{L} \int_0^L x \sin^2\left(\frac{\pi x}{L}\right) dx$$

$$\bar{x} = \frac{L}{2} \quad \overline{x^2} = \int_0^L x^2 \psi^* \psi dx = \frac{2}{L} \int_0^L x^2 \sin^2\left(\frac{\pi x}{L}\right) dx$$

$$\overline{x^2} = \frac{1}{6} \left(2 - \frac{3}{\pi^2}\right) L^2$$

$$\sigma^2 = \overline{x^2} - \bar{x}^2 = \left(\frac{1}{3} - \frac{1}{2\pi^2}\right) L^2 - \frac{L^2}{4} = L^2 \left[\frac{1}{12} - \frac{1}{2\pi^2}\right] = 0.03267 L^2$$

$$\sigma = 0.18076 L$$

Correspondence Principle $n \rightarrow \text{large}$

$$\overline{x^2} \quad \langle x^2 \rangle = L^2 \left(\frac{1}{3} - \frac{1}{2\pi^2 n^2} \right) \quad \bar{x}^2 = \frac{L^2}{4} \quad \sigma^2 = \overline{x^2} - \bar{x}^2 = L^2 \left(\frac{1}{12} - \frac{1}{2\pi^2 n^2} \right)$$

$\psi^* \psi$ is uniformly distributed between $0 \rightarrow L$ for large n .

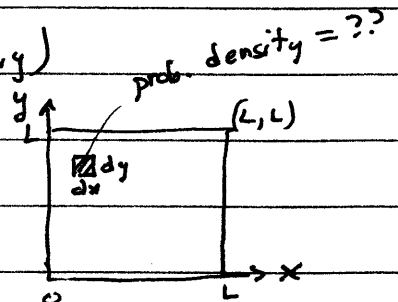
Two Dimensional Infinite Potential Well. (Degeneracy)

start with the Schrödinger Equation. $\rightarrow \psi(x,y)$

$$-\frac{\hbar^2}{2m} \left(\frac{d^2 \psi}{dx^2} + \frac{d^2 \psi}{dy^2} \right) + U(x,y) \psi(x,y) = E \psi(x,y)$$

$$U(x,y) = 0 \quad 0 \leq x \leq L \quad 0 \leq y \leq L$$

$$U(x,y) = \infty \quad \text{otherwise}$$



Assume the solutions are separable.